# Experimental and Theoretical Investigation of the Molecular and Electronic Structure of 3-Acetoxy-2-Methylbenzoic Acid using Quantum Chemical Computational Calculations 

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#### Abstract

In this work, experimental and theoretical study on the molecular structure of 3-acetoxy-2-methyl benzoic acid (3A2MBA) is performed based on the quantum mechanical approach by density functional theory (DFT) calculations. The optimized geometrical parameters were done by DFT using the B3LYP method show good agreement with experimental $X$-ray data. The vibrational harmonic frequencies were scaled using scale factors, yielding a good agreement between the experimentally recorded and the theoretically calculated values with the aid of a calculation of the Total energy distribution (TED). The population of electrons in core, valance and Rydberg sub-shells have been predicted. The natural hybrid atomic orbital studies enhance us to know about the type of orbitals and its percentage of s-type and p-type character. The intra molecular charge transfers occurring in the molecule have been analyzed by NBO analysis. The study is extended to HOMOLUMO analysis, ionization potential, electron affinity, electrophilicity index, chemical potential, electro negativity, softness, hardness and thermodynamic properties of 3A2MBA.


Keywords: 3A2MBA; DFT; FT-IR; FT-Raman; HOMO; LUMO

## INTRODUCTION

Benzoic acid is a compound that has an elegant simplicity to its molecular structure, but its derivatives display an enormous complexity and diversity of molecular structures. The derivatives of benzoic acid have been the subject of investigation as this form an essential component of the vitamin B-complex. Benzoic acid occurs widely in plants and animal tissues along with vitamin B-complex and is used to raise the salicylate level in blood. It is used as antifungal agents for superficial fungus infection of skin. Together with salicylic acids in Whitefield's ointments, which is used for the treatment of fungal skin diseases such as tinea, ring worm in dogs and other species[1]. Benzoic acid and its sodium salt (sodium benzoate) are the most common, safe, food preservatives and antimicrobial agents. Recognizing the commercial availability and low cost of benzoic acid and sodium benzoate, they are found to be attractive candidates to be incorporated into coatings as environmentally benign alternatives [2,3]. The crystal structures of unsubstituted benzoic acid using X-ray diffraction $[4,5]$ as well as the crystal structures of benzoic acid derivatives bearing a halogen atom in an ortho position to the carboxylic acid group are apparent in the literature [6-8]. Crystal structures of 2-acetoxy-3methylbenzoic acid (3-methyl aspirin) and 2-acetoxy-6-methylbenzoic acid (6-methyl aspirin) have already been reported in the literature [9,10]. Aspirin is a unique drug as it is effective against pain, it has anti-pyretic and anti-inflammatory properties, and it is widely used during heart attacks or strokes. A series of water-soluble (benzoyloxy) methyl esters of acetylsalicylic acid (ASA) as pro drugs was synthesized and evaluated by Rolando et al. [11]. Ulla Derthasching et al. have investigated the effects of aspirin on platelet function and
coagulation in human endotoxemia [12]. Investigations on the structure and fundamental vibrations of benzoic acid and its derivatives are still being carried out, increasingly. Vibrational spectra of benzoic acid and substituted benzoic acids have been studied by various workers [13-18]. Amalanathan et al. studied the infrared absorption spectra of 3,5 -dinitrobenzoic acid have been recorded in the solid phase [19]. The stokes and anti stokes laser Raman spectra of 2,3,5-triiodobenzoic acid have been studied [15]. The molecule 3-acetoxy-2methylbenzoic acid is a light yellow crystalline solid; it is used as an intermediate in the preparation of smallsized HIV protease inhibitors [20]. For this reason, we are interested in this class of compounds as potential agents in other diseases. The theoretical calculation of the title molecule is described herein. Previously we have reported crystal structure of 3-acetoxy-2-methylbenzoic acid [21]. Literature survey reveals that to the best of our knowledge no vibrational spectroscopic analysis of 3-acetoxy-2- methyl benzoic acid by the quantum mechanical DFT method have been reported so far. Therefore, the present investigation was undertaken to study the vibrational spectra of this molecule completely and to identify the various normal modes with greater wavenumber accuracy. The theoretical Density Functional Theory (DFT) studies give information regarding the structural parameters, the functional groups, orbital interactions and vibrational frequencies. DFT have evolved to be a powerful and reliable tool for the determination of various molecular properties. B3LYP functional have provided an excellent compromise between accuracy and computational efficiency of vibrational spectra for large and medium size molecule [22-25]. For this reason, by employing the DFT (B3LYP) method, we have calculated the geometric parameters and vibrational spectra of 3-acetoxy-2-methylbenzoic acid in the ground state and compared with the experimental vibrational frequencies. Vibrational assignments were carried out on the basis of total energy distributions (TED) and experimental data. In addition, the ionization potential, electron affinity, electrophilicity index, chemical potential, electronegativity, softness, hardness, HOMO - LUMO energy, natural population analysis, natural orbital analysis and thermodynamic properties calculations were carried out calculations were carried out using B3LYP with $6-31+G(d, p)$ and $6-311++G(d, p)$ method. These calculations were expected to provide new insight in to the vibrational spectra and molecular parameters.

## EXPERIMENTAL SECTION

The FT-IR spectrum of the title molecule was recorded in the region $4000-400 \mathrm{~cm}^{-1}$ on a BRUKER IFS 66 V spectrophotometer using the KBr pellet technique. The spectrum was recorded at room temperature, with a scanning speed of $10 \mathrm{~cm}^{-1}$ per minute and at the spectral resolution of $2.0 \mathrm{~cm}^{-1}$. The FT-Raman spectrum of this molecule was also recorded in the region $3500-50 \mathrm{~cm}^{-1}$ with BRUKER RFS 27 Raman module equipped with Nd:YAG laser source operating at 1064 nm line width 100 mW power. The spectrum was recorded with scanning speed of $50 \mathrm{~cm}^{-1} \mathrm{~min}^{-1}$ of spectral width $4 \mathrm{~cm}^{-1}$. The reported wave numbers are believed to be accurate within $\pm 1 \mathrm{~cm}^{-1}$.

## Computational Methods

Quantum chemical density functional theory calculations were carried out with the 2009 version of the Gaussian program package [26] using B3LYP functions $[27,28]$ combined with the standard $6-31+G(d, p)$ and $6-$ $311++G(d, p)$ basis sets (referred to as small and large basis sets, respectively). The Cartesian representation of the theoretical force constants has been computed at the optimized geometry by assuming $\mathrm{C}_{\mathrm{s}}$ point group symmetry. The optimized geometrical parameters, fundamental vibrational frequencies, IR intensity, Raman activity, reduced mass, force constant and thermodynamic parameters were calculated. From the intensity theory of Raman scattering [29] the relative Raman intensities (Ii) and Raman activities (Si) were calculated:

$$
\begin{equation*}
I_{i}=\frac{\mathrm{f}\left(v_{0}-v_{\mathrm{i}}\right)^{4} \mathrm{~S}_{\mathrm{i}}}{v_{\mathrm{i}}\left(1-\exp \left(-\frac{\mathrm{hc} v_{\mathrm{i}}}{\mathrm{kT}}\right)\right)} \tag{1}
\end{equation*}
$$

Where, $v_{0}$ is the exciting wavenumber; $v_{i}$ the vibrational wavenumber of the normal mode; $h, c$ and $k$ are the universal constants, and $f$ is the suitably chosen common normalization factor for all the peak intensities. Scaling of the force field was performed according to the SQM procedure [30,31] using selective scaling in the natural internal coordinate representation [32]. Transformation of the force field and the subsequent normal coordinate analysis (NCA) including the least squares refinement of the scaling factors, calculation of total energy distribution (TED) were done on a PC with the MOLVIB program (version 7.0-G77) written by Sundius [33-35].

## RESULTS AND DISCUSSION

## Geometrical Parameters

The molecular structure of title molecule belongs to $\mathrm{C}_{\mathrm{s}}$ point group symmetry. The optimized molecular structure of title molecule is obtained from Gaussian 09W as shown in Figure 1. The most optimized structural
parameters of 3-acetoxy-2-methylbenzoic acid are calculated by B3LYP level with $6-31+G(d, p)$ and 6$311++G(d, p)$ basis sets and presented in Table 1. The results are compared with the experimental X-ray diffraction data [21]. From the single crystal XRD data, the title molecule has the cell dimensions are $a=7.754$ $\AA ; \mathrm{b}=11.346 \AA ; \mathrm{c}=21.187 \AA ; \alpha=\gamma=90^{\circ} ; \beta=95^{\circ}$ have been already reported [21]. The values of all the bond lengths and bond angles have small deviations with the experimental (XRD) and theoretical results (which are calculated by B3LYP/6-31+G(d,p) and 6-311++G(d,p) basis sets). This is due to the theoretical calculations are performed (for an isolated molecule) in gaseous phase whereas the experimental results are obtained from solid state. From the crystal data of title molecule, the C11-H12, C11-H13, C11-H14, C18-H19, C18-H20, C18-H21, C4-H22, C5-H23 and C6-H24 bond lengths are shorter than those of the calculated B3LYP values and this is caused by the inter molecular hydrogen bonding in crystalline state. The calculated and experimental C-O group bond distances are almost equal in $\mathrm{C} 7-\mathrm{O} 8, \mathrm{C} 7-\mathrm{O} 9, \mathrm{C} 3-\mathrm{O} 15, \mathrm{C} 16-\mathrm{O} 15$ and $\mathrm{O} 17-\mathrm{C} 16$. Electron donating groups are generally ortho/para directors which electron withdrawing groups are meta directors with the exception of halogens ( $\mathrm{X}=\mathrm{Cl}, \mathrm{F}, \mathrm{I}, \mathrm{Br}$ ). When a functional -COOH group bonded to a benzene ring, the bond angle around C 1 atom (C6-C1-C2) is $120.8^{\circ}$ which confirms the tetrahedral arrangement of C 1 atom. With the electron donating substituents ( $-\mathrm{CH}_{3}$ group) on the benzene ring and the bond angle smaller than $120^{\circ}(\mathrm{C} 1-\mathrm{C} 2-\mathrm{C} 3)$ at the point of substitution bond angle larger than $120^{\circ}$ at the meta position where $-\mathrm{OCOCH}_{3}$ is attached $(\mathrm{C} 4-\mathrm{C} 3-\mathrm{C} 2=$ $128.8^{\circ}$ ). Due to the electron withdrawing nature of the substituents, it is observed that in title molecule the bond angle at the point of substitution C2-C3-O15 is equal values $119.30^{\circ}$ and $119.84^{\circ}$ in B3LYP while the bond angle is $118.6^{\circ}$ by experiment, as shown in Table 1. The calculated bond angles using both basis sets for C2-C3$\mathrm{C} 4, \mathrm{C} 3-\mathrm{C} 4-\mathrm{C} 5, \mathrm{C} 4-\mathrm{C} 5-\mathrm{C} 6, \mathrm{C} 4-\mathrm{C} 5-\mathrm{H} 15, \mathrm{C} 6-\mathrm{C} 5-\mathrm{H} 15, \mathrm{C} 5-\mathrm{C} 6-\mathrm{C} 7, \mathrm{C} 5-\mathrm{C} 6-\mathrm{H} 16$ and C7-C6-H16 are in good agreement with the experimental ones and the small variation in remaining bond angles depends on the electro negativity of the central atom, the presence of lone pair of electrons and the conjugation of the double bonds.


Figure 1: Molecular structure and atomic labelling of 3A2MBA

## Vibrational Spectra

In our present study, we have performed a frequency calculation analysis to obtain the spectroscopic signature of 3-acetoxy-2-methylbenzoic acid (3A2MBA). The molecule 3A2MBA has an orthorhombic structure of Pbca space group and Cs point group symmetry. The vibrational spectral analysis of the molecule 3A2MBA is done based on normal coordinate analysis followed by scaled quantum mechanical force field calculations. The 3A2MBA molecule consists of 24 atoms therefore they have 66 vibrational normal modes. Out of 66 modes of fundamental vibrations 45 modes should be in-plane symmetric ( $A^{\prime}$ ) and 21 modes out-of-plane symmetric ( $A^{\prime \prime}$ ) with respect to the reflection on the symmetry plane. All the frequencies are assigned in terms of fundamental, overtone and combination bands. The recorded (FT-IR and FT-Raman) and calculated vibrational wavenumbers along with their relative intensities and probable assignments with TED of 3A2MBA molecule are given in Table 2. The theoretical spectra were obtained from the B3LYP/6-31+G ( $\mathrm{d}, \mathrm{p}$ ) and $6-311++\mathrm{G}(\mathrm{d}, \mathrm{p})$ method using Lorentzian band shape with band width on half-height $10 \mathrm{~cm}^{-1}$. This reveals good correspondence between theory and experiment in main spectral features. The experimental spectra are shown in Figures 2 and 3, respectively. The computed vibrational wavenumbers and the atomic displacements corresponding to the different normal modes are used for identifying the vibrational modes unambiguously. In order to simulate H bonding through COOH group we also calculated the vibrational wavenumbers of 3 A 2 MBA molecule. The calculated wavenumbers are usually higher than the corresponding experimental quantities, due to the combination of electron correlation effects and basis sets deficiencies. After applying, the different scaling factors, the theoretical wavenumbers are in good agreement with experimental wavenumbers [26].


Figure 2: FT-IR spectrum of 3A2MBA


Figure 3: FT-Raman spectrum of 3A2MBA

## Carboxylic Acid (-COOH) Group Vibrations

The vibrational analysis of carboxylic acid $(-\mathrm{COOH})$ group is made on the basis of carbonyl group and hydroxyl group vibrations. The carboxylic group is substituted at the first position in benzene ring of the 3A2MBA molecule. The vibrational analysis of - COOH group is significant because of the activity of the title molecule which is mainly due to either the presence of this moiety or a group that is easily converted to it within the plant tissues [36] and which is made on the basis of carbonyl group and the hydroxyl group. The vibrational bands of the -COOH groups of 3 A 2 MBA contain the $\mathrm{C}-\mathrm{O}, \mathrm{C}=\mathrm{O}$ and $\mathrm{O}-\mathrm{H}$ vibrational modes. $\mathrm{C}=\mathrm{O}$ stretching band appears strongly in the region $1870-1540 \mathrm{~cm}^{-1}$ in which the position of $\mathrm{C}=\mathrm{O}$ stretching band depends on the physical state, electronic and mass effects of neighboring substituents, conjugations, intra molecular and intermolecular hydrogen bonding [37-44]. The O-H group stretching vibrations are likely to be most sensitive to the environment, so they show pronounced shifts in the spectra of the hydrogen bonded species. A free hydroxyl group or a non-hydrogen bonded hydroxyl group absorbs in the range $3700-3500 \mathrm{~cm}^{-1}$, whereas the existence of intermolecular hydrogen bond formation can lower the $\mathrm{O}-\mathrm{H}$ stretching frequency to the $3550-3200 \mathrm{~cm}^{-1}$ region with increase in intensity and breath [39-42]. In the present investigation, the stretching vibration of $\mathrm{O}-\mathrm{H}$ group is assigned as a very weak band observed at $3317 \mathrm{~cm}^{-1}$ in the recorded FT-IR spectrum. The $\mathrm{O}-\mathrm{H}$ stretching modes are found as 3329 and $3322 \mathrm{~cm}^{-1}$ by B3LYP calculation levels, as given in Table 2. The theoretically computed scaled frequency for $\mathrm{O}-\mathrm{H}$ vibration by B3LYP/6-311++G(d,p) method shows excellent agreement with recorded spectrum as well as the literature data and also find support from TED values of $100 \%$. The lower stretching frequency observed in 3A2MBA compared with the free $\mathrm{O}-\mathrm{H}$ group stretching signifies that there is a possibility of intermolecular hydrogen bonding in 3A2MBA, between the hydroxyl group and carbonyl group of molecule. The $\mathrm{O}-\mathrm{H}$ in-plane bending vibrations occur in the general region of $1420-1330 \mathrm{~cm}^{-1}$ and are not much affected due to hydrogen bonding unlike the stretching and out-of-plane bending frequencies. The in-plane banding mode of hydroxyl group was recorded at $1375 \mathrm{~cm}^{-1}$ in FT-Raman spectrum and computed at 1384 and $1381 \mathrm{~cm}^{-1}$ with $53 \%$ value of TED contribution for title molecule. They show good agreement with the theoretically computed B3LYP/6-311++G (d,p) method. The O-H out-of-plane bending vibrations appear in the region $320-290 \mathrm{~cm}^{-1}$ for free $\mathrm{O}-\mathrm{H}$ and in the region $710-570 \mathrm{~cm}^{-1}$ for associated benzoic acids [43]. The carboxylic acids show out-of-plane bending band of $\mathrm{O}-\mathrm{H}$ vibration occurred at $704 \mathrm{~cm}^{-1}$ in FT-Raman spectrum of 3A2MBA. The TED calculations show that the hydroxyl stretching vibrational mode is very pure. But the inplane bending vibration of the hydroxyl group is overlapped with the other vibrations. The band due to the free hydroxyl group is sharp and its intensity increases. The $\mathrm{C}=\mathrm{O}$ stretching mode is the strongest bands in the infrared and appears with diminished intensity weak bands in the Raman spectra (less polarizability resulting due to highly dipolar carbonyl bond) around $1800-1650 \mathrm{~cm}^{-1}$ in aromatic compounds is the most salient feature of the presence of carbonyl group and are due to the $\mathrm{C}=\mathrm{O}$ stretching motions [43,44]. The wavenumber of the $\mathrm{C}=\mathrm{O}$ stretch due to carbonyl group mainly depends on the bond strength, which in turn depends upon inductive,
conjugative, steric effects and the lone pair of electrons on oxygen. Hence the weak FT-Raman band observed at $1769 \mathrm{~cm}^{-1}$ is assigned to the $\mathrm{C}=\mathrm{O}$ stretching band of 3A2MBA molecule. As indicated by TED, this mode involves exact contribution of $89 \%$ suggesting that it is a pure stretching mode. The similar vibrations are calculated at 1781 and $1778 \mathrm{~cm}^{-1}$ by DFT-B3LYP/6-31+G (d,p) and 6-311++G (d,p) method. However, the computed scaled values observed at 1776 and $1773 \mathrm{~cm}^{-1}$ of $\mathrm{C}=\mathrm{O}$ stretching vibrations shows better agreement with the experimental data. However, it is known that $\mathrm{C}=\mathrm{O}$ stretching and $\mathrm{O}-\mathrm{H}$ bending modes are not independent vibrational modes because they coupled with the vibrations of adjacent groups. The $\mathrm{C}-\mathrm{O}$ stretching band is occurs at $1450-1150 \mathrm{~cm}^{-1}$ regions depending on whether monomer or other hydrogen bonded species are present. Generally the $\mathrm{C}-\mathrm{O}$ stretching mode appears at lower frequency than the $\mathrm{C}-\mathrm{OH}$ bending mode. However, these bands overlap with other bands that are due to aromatic or aliphatic chain vibrations. Based on the above information, the FT-Raman band observed at $1262 \mathrm{~cm}^{-1}$ has been assigned to $\mathrm{C}-\mathrm{O}$ stretching vibration for 3 A 2 MBA . In addition carboxylic acids also show $\mathrm{C}=\mathrm{O}$ and $\mathrm{C}-\mathrm{O}$ in-plane and out-of-plane bending vibrations are also listed in Table 2.

## C-H Vibrations

In general, the aromatic structure shows presence of C-H stretching vibration in the band region 3100-3000 $\mathrm{cm}^{-1}$ [45]. Bayari et al. have reported $\mathrm{C}-\mathrm{H}$ stretching frequencies in the range of $3125-3000 \mathrm{~cm}^{-1}$ [46]. In the present investigation, three weak bands observed in FT-Raman spectrum at 3125,3100 and $3075 \mathrm{~cm}^{-1}$ are assigned to $\mathrm{C} 4-\mathrm{H} 22, \mathrm{C} 5-\mathrm{H} 23$ and $\mathrm{C} 6-\mathrm{H} 24$ stretching vibrations. The theoretically calculated wave numbers agree closely with the experimental spectral values as seen in Table 2. These assignments are well supported by calculated TED values. All bands have weak intensities and were obtained in the expected region. As expected, these three modes are pure stretching modes as it is evident from TED column. The $\mathrm{C}-\mathrm{H}$ in-plane bending vibrations are assigned in the region $1300-1000 \mathrm{~cm}^{-1}$ [35]. For this compound, the $\mathrm{C}-\mathrm{H}$ in-plane bending vibrations were observed at 1070,1023 and $1011 \mathrm{~cm}^{-1}$ as strong and medium strong in FT-IR spectrum. The theoretically scaled vibrations by B3LYP/6-31+G(d,p) and $6-311++G(d, p)$ levels also shows good agreement with the experimentally recorded data. The C-H out-of-plane bending vibration appears within the region 1000$700 \mathrm{~cm}^{-1}$ [35]. The weak vibrations were identified at 776,764 in FT-IR and $781 \mathrm{~cm}^{-1}$ in FT-Raman, and they are assigned to C-H out-of-plane bending vibrations for title molecule. These modes were computed at 803, 777, $730 \mathrm{~cm}^{-1}$ (unscaled)/ $766,735,713 \mathrm{~cm}^{-1}$ (scaled) and 801, $732,778 \mathrm{~cm}^{-1}$ (unscaled) /764, 736, $711 \mathrm{~cm}^{-1}$ (scaled) by the B3LYP with $6-31+G(d, p)$ and $6-311++G(d, p)$ method. The C-H out-of-plane bending vibrations are also lie within the characteristic region and after scaling procedure, the theoretically scaled $\mathrm{C}-\mathrm{H}$ vibrations are found to be in good agreement with the experimental values and literature [47]. The change in the frequencies of these deformations from the values in 3A2MBA is determined mainly by the relative position of the substituent and is almost independent of their nature.

## C-C Vibrations

Vibrations of phenyl ring have been comprehensively studied according to Wilson's numbering convention [48]. The phenyl ring vibrational modes of title molecule have been analyzed based on the vibrational spectra of previously published vibrations of the benzene molecule is helpful in the identification of the phenyl ring modes [49-52]. The $\mathrm{C}-\mathrm{C}$ ring stretching vibrations are very prominent, as the double bond is in conjugation with the ring, in the vibrational spectra of benzene and its derivatives [48]. The $\mathrm{C}-\mathrm{C}$ stretching vibrations of the phenyl group are expected in the range $1650-1200 \mathrm{~cm}^{-1}$ [48]. In general, the five bands of variable intensity observed at 1625-1590, 1590-1575, 1540-1470, 1465-1430 and 1380-1280 $\mathrm{cm}^{-1}$ have been already reported by Varsanyi [44]. In the present work, the FT-IR and FT-Raman spectra bands of variable intensity observed at 1623(w), $1588(\mathrm{w}), 1482(\mathrm{~s}), 1129(\mathrm{~s}) \mathrm{cm}^{-1}$ and 1638(mw), 1602(ms), 1576(mw), 1087(s), 807(w) $\mathrm{cm}^{-1}$ have been assigned to $\mathrm{C}-\mathrm{C}$ stretching vibrations in 3A2MBA. The theoretically computed (scaled) values at $1645,1627,1608$, 1592, 1581, 1487, 1134, 1092 and $771 \mathrm{~cm}^{-1}$ by the B3LYP/6-311++G (d,p) method show an excellent agreement with experimental values. The $\mathrm{C}-\mathrm{C}-\mathrm{C}$ in-plane and out-of-plane bending vibrations are also calculated and the correlation between the experimental and theoretical values are observed as shown in Table 2. In lower frequency region $\mathrm{C}-\mathrm{C}$ stretching and bending vibrations modes mixed with other vibrational modes and with small amount contributions were assigned based on TED results. From TED results, the most of these modes were assigned as pure vibrations. All the bands lie in the expected range to literature values. These observed frequencies show that, the substitutions in the ring to some extend affect the ring modes of vibrations. The comparison of the theoretically computed values were very good in agreement by B3LYP/6-311++G (d,p) method.

Table 1: Calculated optimized geometrical parameters of 3A2MBA

| Bond lengths ( $\mathbf{( 1 )}$ | Value ( ${ }_{\text {A }}$ ) |  |  | Bond angles (o) | Value ( ${ }^{\circ}$ ) |  |  | Dihedral angles (o) | Value ( ${ }^{\circ}$ ) |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ${ }^{\text {a }}$ Expt. | B3LYP |  |  | ${ }^{\text {a }}$ Expt. | B3LYP |  |  | ${ }^{\text {a }}$ Expt. | B3LYP |  |
|  |  | 6-31+G(d,p) | 6-311++G(d,p) |  |  | 6-31+G(d,p) | 6-311++G(d,p) |  |  | 6-31+G(d,p) | 6-311++G(d,p) |
| C1-C2 | 1.397 | 1.419 | 1.415 | C2-C1-C6 | 120 | 120.84 | 120.42 | C6-C1-C2-C3 | 0.1 | -2 | -2.1 |
| C1-C6 | 1.384 | 1.405 | 1.401 | C2-C1-C7 | 122.1 | 124.97 | 124.85 | C6-C1-C2-C11 | -179.8 | 176.6 | 176.1 |
| C1-C7 | 1.462 | 1.496 | 1.496 | C6-C1-C7 | 118 | 114.2 | 114.33 | C7-C1-C2-C3 | 179.7 | 178.1 | 177.9 |
| C2-C3 | 1.387 | 1.407 | 1.404 | C1-C2-C3 | 116.2 | 116.2 | 116.31 | C7-C1-C2-C11 | -0.2 | -3.2 | -3.7 |
| C2-C11 | 1.49 | 1.512 | 1.51 | C1-C2-C11 | 124.4 | 124.72 | 124.62 | C2-C1-C6-C5 | 0.4 | 0.6 | 0.8 |
| C3-C4 | 1.374 | 1.394 | 1.391 | C3-C2-C11 | 119.5 | 119.07 | 119.04 | C2-C1-C6-H24 | - | -179 | -178.7 |
| C3-O15 | 1.397 | 1.397 | 1.395 | C2-C3-C4 | 123.6 | 123.04 | 122.34 | C7-C1-C6-C5 | -179.3 | -179.4 | -179.2 |
| C4-C5 | 1.363 | 1.393 | 1.39 | C2-C3-015 | 118.6 | 119.3 | 119.84 | C7-C1-C6-H24 | - | 0.8 | 1 |
| C4-H22 | 0.93 | 1.086 | 1.084 | C4-C3-015 | 117.6 | 117.46 | 117.64 | C2-C1-C7-O8 | -12 | 167.9 | 163.2 |
| C5-C6 | 1.364 | 1.391 | 1.387 | C3-C4-C5 | 118.9 | 119.6 | 119.66 | C2-C1-C7-09 | 169 | -13.6 | -18.6 |
| C5-H23 | 0.93 | 1.085 | 1.083 | C3-C4-H22 | 121 | 118.2 | 118.94 | C6-C1-C7-08 | 167.6 | -11.9 | -16.5 |
| C6-H24 | 0.93 | 1.084 | 1.082 | C5-C4-H22 | 121 | 121.46 | 121.39 | C6-C1-C7-09 | -10.6 | 166.5 | 161.4 |
| C7-08 | 1.228 | 1.217 | 1.209 | C4-C5-C6 | 119.6 | 119.25 | 119.26 | C1-C2-C3-C4 | -0.7 | 2.5 | 2.5 |
| C7-09 | 1.3 | 1.357 | 1.356 | C4-C5-H23 | 120 | 120.34 | 120.32 | C1-C2-C3-015 | 175 | 177.3 | 177.4 |
| O9-H10 | 0.93 | 0.973 | 0.969 | C6-C5-H23 | 120 | 120.41 | 120.41 | C11-C2-C3-C4 | 179.2 | -176.1 | -175.8 |
| C11-H12 | 0.96 | 1.091 | 1.089 | C1-C6-C5 | 121.8 | 121.04 | 121 | C11-C2-C3-O15 | -5.1 | -1.39 | -1 |
| C11-H13 | 0.96 | 1.095 | 1.088 | C1-C6-H24 | 119 | 118.03 | 118.12 | C1-C2-C11-H12 | - | -45.1 | -42.2 |
| C11-H14 | 0.96 | 1.09 | 1.093 | C5-C6-H24 | 119 | 120.93 | 120.87 | C1-C2-C11-H13 | - | 73.4 | 76.3 |
| O15-C16 | 1.397 | 1.388 | 1.38 | C1-C7-08 | 121.3 | 123.96 | 124.09 | C1-C2-C11-H14 | - | -165.7 | -162.8 |
| C16-O17 | 1.1182 | 1.205 | 1.197 | C1-C7-09 | 114.8 | 115.1 | 114.52 | C3-C2-C11-H12 | - | 133.4 | 135.9 |
| C16-C18 | 1.469 | 1.512 | 1.51 | O8-C7-H9 | - | 121.01 | 121.35 | C3-C2-C11-H13 | - | -108 | -105.4 |
| C18-H19 | 0.96 | 1.094 | 1.092 | C7-O9-H10 | 112 | 106.07 | 106.34 | C3-C2-C11-H14 | - | 12.8 | 15.3 |
| C18-H20 | 0.96 | 1.094 | 1.092 | C2-C11-H12 | 109 | 111.19 | 111.23 | C2-C3-C4-C5 | 0.9 | -1.7 | -1.6 |
| C18-H21 | 0.96 | 1.089 | 1.088 | C2-C11-H13 | 109 | 111.46 | 111.4 | C2-C3-C4-H22 | - | 177.4 | 177.4 |
|  |  |  |  | C2-C11-H14 | 109 | 110.3 | 110.24 | O15-C3-C4-C5 | -174.9 | -176.5 | -176.4 |
|  |  |  |  | H12-C11-H13 | 110 | 106.42 | 106.49 | O15-C3-C4-H22 | - | 2.6 | 2.6 |
|  |  |  |  | H12-C11-H14 | 109 | 108.61 | 108.66 | C2-C3-O15-C16 | - | 94.9 | 95.6 |
|  |  |  |  | H13-C11-H14 | 110 | 108.73 | 108.66 | C4-C3-O15-C16 | - | -89.9 | -89.3 |
|  |  |  |  | C3-015-C16 | 119 | 122.31 | 122.35 | C3-C4-C5-C6 | -0.5 | 0.2 | 0.2 |
|  |  |  |  | O15-C16-017 | 121.6 | 117.86 | 117.93 | C3-C4-C5-H23 | - | 179.9 | 179.8 |
|  |  |  |  | O15-C16-C18 | 112.1 | 117.56 | 117.3 | H22-C4-C5-C6 | - | -178.9 | -178.8 |
|  |  |  |  | 017-C16-C18 | 126.4 | 124.58 | 124.75 | H22-C4-C5-H23 | - | 0.7 | 0.7 |
|  |  |  |  | C16-C18-H19 | 109 | 111.15 | 111.14 | C4-C5-C6-C1 | -0.2 | 0.2 | 0.1 |
|  |  |  |  | C16-C18-H20 | 109 | 110.97 | 110.83 | C4-C5-C6-H24 | - | 179.9 | 179.7 |
|  |  |  |  | C16-C18-H21 | 109 | 108.08 | 108.13 | H23-C5-C6-H24 | - | 0.2 | 0.1 |
|  |  |  |  | H19-C18-H20 | 110 | 107.22 | 107.23 | C1-C7-09-O10 | - | -179 | -179.1 |
|  |  |  |  | H19-C18-H21 | 110 | 109.83 | 109.89 | O8-C7-O9-H10 | 0.8 | 0.5 | -0.9 |



Table 2: Experimental and calculated (FT-IR and FT-Raman) frequencies (cm-1), IR intensity (KM mol-1), Raman activity (Å4 amu-1), force constants (m dyne $\AA$-1), reduced mass (amu) and assignments for 3A2MBA

| S No. | Experimental frequencies |  | Calculated with B3LYP/6-31+G(d,p) |  |  |  |  |  | Calculated with B3LYP/6-311++G(d,p) |  |  |  |  |  | Vibrational Assignments along with \% of TED |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | FT-IR | FT-Raman | Un | Scaled | ${ }^{\text {a }}$ IR int. | ${ }^{\text {b Ra. act. }}$ | Force | Reduced | Un | Scaled | ${ }^{\text {a }}$ IR int. | ${ }^{\text {b Ra. act. }}$ | Force | Reduced |  |
|  |  |  | scaled |  |  |  | constants | mass | scaled |  |  |  | constants | mass |  |
| 1. | 3317 vw |  | 3752 | 3329 | 96.0305 | 151.12 | 8.8302 | 1.0643 | 3756 | 3322 | 99.8576 | 157.14 | 8.8508 | 1.0644 | vO-H(100) |
| 2. |  | 3125w | 3229 | 3136 | 2.8879 | 121.59 | 6.7212 | 1.0939 | 3209 | 3132 | 2.8469 | 123.19 | 6.6419 | 1.0942 | vC-H (98) |
| 3. |  | 3100w | 3212 | 3108 | 4.0237 | 141.16 | 6.6531 | 1.0945 | 3193 | 3105 | 3.2361 | 143.34 | 6.5753 | 1.094 | vC-H(98) |
| 4. |  | 3075w | 3198 | 3084 | 4.2768 | 63.26 | 6.5555 | 1.0879 | 3179 | 3081 | 3.4872 | 72.9 | 6.4804 | 1.0877 | vC-H(99) |
| 5. |  | 2999w | 3178 | 3006 | 6.091 | 76.73 | 6.5669 | 1.1032 | 3158 | 3002 | 7.2377 | 76.73 | 6.4833 | 1.1027 | $\mathrm{CH}_{3} \mathrm{Ops}(96)$ |
| 6. | 2941w |  | 3164 | 2953 | 5.4306 | 45.09 | 6.5093 | 1.103 | 3145 | 2949 | 5.0973 | 54.78 | 6.4305 | 1.1029 | $\mathrm{CH}_{3} \mathrm{ops}(96)$ |
| 7. |  | 2920w | 3130 | 2928 | 2.8377 | 23.9 | 6.3523 | 1.1 | 3112 | 2926 | 2.9016 | 26.98 | 6.2756 | 1.0994 | $\mathrm{CH}_{3} \mathrm{ss}(95)$ |
| 8. | 2894w |  | 3129 | 2902 | 8.5916 | 74.97 | 6.3037 | 1.0928 | 3111 | 2899 | 9.1109 | 83.78 | 6.2192 | 1.09 | $\mathrm{CH}_{3} \mathrm{ss}(95)$ |
| 9. | 2882w |  | 3065 | 2894 | 2.5159 | 100 | 5.7437 | 1.0372 | 3053 | 2888 | 2.0183 | 98.89 | 5.6953 | 1.0368 | $\mathrm{CH}_{3} \mathrm{ips}(94)$ |
| 10. |  | 2868w | 3062 | 2858 | 10.6401 | 87.22 | 5.7639 | 1.0431 | 3048 | 2854 | 9.366 | 87.22 | 5.7193 | 1.0446 | $\mathrm{CH}_{3} \mathrm{ips}(94)$ |
| 11. | 2847mw |  | 1837 | 1976 | 445.4052 | 19 | 23.345 | 11.7311 | 1833 | 1971 | 457.7305 | 29.05 | 23.2448 | 11.7337 | $\begin{gathered} \text { Rtrigd(56), bC- } \\ \mathrm{C}(23) \\ \hline \end{gathered}$ |
| 12. |  | 1769w | 1781 | 1776 | 410.8519 | 82.68 | 18.1269 | 9.6993 | 1778 | 1773 | 428.2895 | 82.78 | 17.9724 | 9.647 | $v \mathrm{C}=\mathrm{O}(84)$ |
| 13. | 1763s |  | 1643 | 1772 | 24.4671 | 54.84 | 9.9001 | 6.2183 | 1635 | 1769 | 22.315 | 56.8 | 9.7468 | 6.1827 | $\mathrm{vC}=\mathrm{O}(72)$ |


| 14. | 1694vs |  | 1617 | 1702 | 22.9531 | 14.45 | 9.5189 | 6.1739 | 1610 | 1699 | 22.9094 | 14.68 | 9.301 | 6.0871 | $\begin{gathered} \operatorname{Rsymd}(59), \\ \text { Rasymd(29), vC- } \\ \mathrm{C}(19) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 15. |  | 1638mw | 1506 | 1648 | 11.1238 | 5.53 | 1.5182 | 1.1357 | 1503 | 1645 | 10.7164 | 15.54 | 1.5096 | 1.334 | $\begin{gathered} \text { vC-C(71), vC- } \\ \mathrm{H}(16), \mathrm{bC}-\mathrm{H}(10) \end{gathered}$ |
| 16. | 1623w |  | 1502 | 1631 | 38.9715 | 4.63 | 2.0091 | 1.5108 | 1498 | 1627 | 36.3885 | 4.73 | 1.9393 | 1.4661 | $\begin{gathered} v C-C(65), b C- \\ \mathrm{H}(11), \mathrm{vC}-\mathrm{H}(13) \end{gathered}$ |
| 17. |  | 1602 ms | 1486 | 1612 | 7.8827 | 3.84 | 2.3587 | 1.8109 | 1481 | 1608 | 8.7691 | 13.84 | 2.1665 | 1.6757 | $\begin{gathered} \text { vC-C(69),bC- } \\ \mathrm{H}(11), \mathrm{bC}-\mathrm{C}(14) \end{gathered}$ |
| 18. | 1588w |  | 1484 | 1595 | 23.5606 | 5.42 | 1.3876 | 1.0686 | 1480 | 1592 | 20.7156 | 6.42 | 1.4692 | 1.1382 | $\begin{gathered} v C-C(58), \text { bC- } \\ \mathrm{H}(15), \omega \mathrm{C}-\mathrm{O}(15) \end{gathered}$ |
| 19. |  | 1576mw | 1475 | 1583 | 31.6921 | 2.91 | 2.3499 | 1.831 | 1471 | 1581 | 33.556 | 2.67 | 2.1437 | 1.6811 | $\begin{gathered} v \mathrm{c}-\mathrm{C}(59), \omega \mathrm{C}- \\ \mathrm{H}(15), \mathrm{bC}-\mathrm{C}(11) \end{gathered}$ |
| 20. | 1482s |  | 1468 | 1489 | 7.8665 | 6.65 | 1.4472 | 1.1394 | 1464 | 1487 | 7.8848 | 6.69 | 1.5141 | 1.1981 | $\begin{aligned} & \text { vC-C(54),bC- } \\ & \mathrm{H}(11), \text { vC-C(6) } \end{aligned}$ |
| 21. | 1425 ms |  | 1425 | 1432 | 6.3769 | 11.26 | 1.4568 | 1.2173 | 1421 | 1429 | 6.4098 | 11.22 | 1.4493 | 1.218 | $\mathrm{CH}_{3} \mathrm{ipb}$ (89) |
| 22. |  | 1412 ms | 1403 | 1421 | 43.0211 | 1.37 | 1.4947 | 1.2881 | 1397 | 1418 | 38.7472 | 11.38 | 1.467 | 1.2744 | $\mathrm{CH}_{3} \mathrm{ipb}$ (89) |
| 23. |  | 1375w | 1372 | 1384 | 122.7663 | 14.81 | 3.1331 | 2.8217 | 1364 | 1381 | 101.7417 | 14.87 | 2.8145 | 2.5664 | $\begin{gathered} \text { bO-H(53), vC-C } \\ (18), \text { bC-C(13) } \end{gathered}$ |
| 24. | 1364ms |  | 1331 | 1371 | 18.1253 | 6.83 | 6.7414 | 6.4547 | 1316 | 1369 | 13.546 | 16.83 | 6.2122 | 6.0836 | $\mathrm{CH}_{3} \mathrm{sb}(63)$ |
| 25. | 1317s |  | 1263 | 1324 | 70.0042 | 13.05 | 2.1158 | 2.2489 | 1259 | 1322 | 47.043 | 13.15 | 1.9958 | 2.1358 | $\mathrm{CH}_{3} \mathrm{sb}(63)$ |
| 26. |  | 1300w | 1236 | 1305 | 122.5571 | 49.89 | 2.0635 | 2.2897 | 1234 | 1302 | 127.8708 | 47.89 | 2.663 | 2.5231 | $\mathrm{CH}_{3} \mathrm{Opb}(54)$ |
| 27. | 1270 ms |  | 1219 | 1278 | 288.3423 | 7.03 | 2.719 | 3.1035 | 1213 | 1276 | 97.7131 | 17.03 | 2.1724 | 2.5024 | $\mathrm{CH}_{3} \mathrm{Opb}$ (33) |


| 28. |  | 1262w | 1207 | 1272 | 266.3554 | 9.39 | 2.1211 | 2.4704 | 1196 | 1268 | 377.9956 | 8.59 | 2.1041 | 2.4947 | $\begin{gathered} \text { vC-O (46), vC-C } \\ (15), \mathrm{bC}-\mathrm{O}(11) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 29. | 1211vs |  | 1190 | 1219 | 48.8286 | 7.5 | 1.2182 | 1.4577 | 1187 | 1217 | 129.4373 | 7.59 | 1.3292 | 1.5994 | $\begin{gathered} \text { vC-O (46), vC-C } \\ (15), \mathrm{bC}-\mathrm{O}(11) \end{gathered}$ |
| 30. |  | 1200 ms | 1141 | 1208 | 151.6279 | 6.24 | 2.108 | 2.7476 | 1133 | 1204 | 165.2799 | 17.23 | 2.2107 | 2.9199 | $\begin{gathered} \text { vC-O (46), vC-C } \\ (15), \mathrm{bC}-\mathrm{O}(11) \end{gathered}$ |
| 31. | 1129s |  | 1102 | 1136 | 30.7544 | 12.95 | 1.4755 | 2.0619 | 1099 | 1134 | 34.8299 | 11.95 | 1.4816 | 2.0784 | $\begin{gathered} v \mathrm{C}-\mathrm{C}(59), \omega \mathrm{C}- \\ \mathrm{H}(15), \mathrm{bC}-\mathrm{C}(11) \end{gathered}$ |
| 32. |  | 1087s | 1062 | 1095 | 18.2063 | 0.66 | 1.1939 | 1.7957 | 1061 | 1092 | 8.9263 | 1.66 | 1.0941 | 1.6484 | $\begin{aligned} & \text { vC-C(54),bC- } \\ & \mathrm{H}(11), \text { vC-C(6) } \end{aligned}$ |
| 33. | 1070ms |  | 1060 | 1079 | 19.6232 | 2.14 | 1.1685 | 1.7648 | 1056 | 1077 | 27.4644 | 3.15 | 1.2587 | 1.9131 | $\begin{aligned} & \text { bC-H (71), bC- } \\ & \text { C(10), vC-O(13) } \end{aligned}$ |
| 34. | 1023s |  | 1056 | 1035 | 9.6063 | 0.24 | 1.1351 | 1.7276 | 1055 | 1031 | 20.8739 | 0.26 | 1.1796 | 1.7967 | $\begin{aligned} & \text { bC-H (33), vC- } \\ & \text { O(13), bC-C(13) } \end{aligned}$ |
| 35. | 1011ms | 996w | 1008 | 1021 | 68.7894 | 1.68 | 1.0724 | 1.7887 | 1005 | 1018 | 63.2646 | 1.78 | 0.9956 | 1.6703 | $\begin{aligned} & \text { bC-H(60), bC- } \\ & \text { C(11), vC-O(10) } \end{aligned}$ |
| 36. | 936s |  | 1000 | 949 | 1.3087 | 0.16 | 0.7852 | 1.3302 | 1002 | 946 | 25.6652 | 0.19 | 0.8714 | 1.4729 | $\mathrm{CH}_{3} \mathrm{ipr}$ (69) |
| 37. | 882mw |  | 952 | 891 | 23.9304 | 2.84 | 1.6782 | 3.1416 | 950 | 889 | 6.1721 | 2.81 | 0.9744 | 1.8309 | $\mathrm{CH}_{3} \mathrm{ipr}(65)$ |
| 38. |  | 871w | 949 | 845 | 5.6575 | 0.66 | 0.7538 | 1.4194 | 948 | 842 | 35.9437 | 1.66 | 1.1378 | 2.1477 | $\mathrm{CH}_{3} \operatorname{opr}(69)$ |
| 39. | 835mw |  | 836 | 786 | 10.5158 | 2.55 | 1.1103 | 2.6942 | 837 | 784 | 7.7372 | 2.59 | 1.0998 | 2.6625 | $\mathrm{CH}_{3} \mathrm{opr}(65)$ |
| 40. |  | 807w | 827 | 773 | 7.5315 | 3.77 | 1.3072 | 3.2387 | 827 | 771 | 7.8878 | 3.37 | 1.4487 | 3.5901 | $\begin{gathered} \text { vC-C(44), bC- } \\ \mathrm{H}(16), \mathrm{bC}-\mathrm{O}(11) \end{gathered}$ |


| 41. |  | 781w | 803 | 766 | 20.6397 | 2.39 | 1.9971 | 5.2506 | 801 | 764 | 28.5687 | 2.99 | 2.0267 | 5.356 | $\begin{aligned} & \omega \mathrm{C}-\mathrm{H}(24), \mathrm{vC}- \\ & \mathrm{O}(10), \omega \mathrm{C}-\mathrm{C}(13) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 42. | 776mw |  | 777 | 735 | 57.5217 | 0.29 | 1.0401 | 2.9172 | 778 | 734 | 51.5039 | 1.29 | 1.0264 | 2.8745 | $\begin{aligned} & \omega \mathrm{C}-\mathrm{H}(26), \mathrm{vC}- \\ & \mathrm{O}(8), \omega \mathrm{C}-\mathrm{C}(10) \end{aligned}$ |
| 43. | 764mw |  | 730 | 713 | 10.4674 | 8.1 | 0.9081 | 2.889 | 732 | 711 | 11.6958 | 8.16 | 0.8378 | 2.6469 | $\begin{aligned} & \omega \mathrm{C}-\mathrm{H}(24), v \mathrm{c}- \\ & \mathrm{O}(16), \omega \mathrm{C}-\mathrm{C}(8) \end{aligned}$ |
| 44. |  | 756mw | 702 | 663 | 9.102 | 13.62 | 1.0862 | 3.7366 | 705 | 660 | 9.9795 | 13.72 | 1.1699 | 3.9882 | Rasymd(61), Rsymd(31), $\omega$ CC(11) |
| 45. | 729w |  | 651 | 632 | 46.5357 | 2.27 | 0.8042 | 3.2138 | 652 | 628 | 42.3604 | 2.37 | 1.0192 | 4.0668 | $\begin{gathered} \text { bC-C(35), bC- } \\ \mathrm{O}(13) \end{gathered}$ |
| 46. |  | 704mw | 627 | 621 | 95.3615 | 3.1 | 0.4099 | 1.7654 | 620 | 619 | 95.2087 | 3.17 | 0.3819 | 1.6824 | $\omega \mathrm{O}-\mathrm{H}(22)$ |
| 47. | 654vw | 615w | 585 | 596 | 9.0229 | 7.35 | 1.0348 | 5.1153 | 586 | 594 | 9.4727 | 8.35 | 1.0057 | 4.968 | $\begin{gathered} \text { bC-C(17), bO- } \\ \mathrm{C}(21), \mathrm{bC}-\mathrm{O}(14), \end{gathered}$ |
| 48. | 623w |  | 567 | 574 | 8.1721 | 1.77 | 0.812 | 4.2813 | 570 | 571 | 9.5917 | 3.73 | 0.7044 | 3.6698 | $\begin{gathered} \text { bC-C(13), bO- } \\ \text { C(42) } \end{gathered}$ |
| 49. |  | 589w | 564 | 553 | 1.9819 | 0.58 | 0.5152 | 2.744 | 566 | 549 | 1.8823 | 0.56 | 0.5135 | 2.7196 | $\begin{gathered} \mathrm{bC}=\mathrm{O}(13), \omega \mathrm{C}- \\ \mathrm{C}(11) \end{gathered}$ |
| 50. |  | 564w | 558 | 549 | 3.8108 | 2.26 | 0.6009 | 3.2681 | 555 | 546 | 5.1211 | 2.28 | 0.5871 | 3.2278 | $\begin{gathered} \mathrm{bC}=\mathrm{O}(32), \mathrm{bC}- \\ \mathrm{C}(10) \end{gathered}$ |
| 51. | 541w |  | 497 | 537 | 0.5348 | 0.63 | 0.4658 | 3.1962 | 495 | 534 | 1.0079 | 0.53 | 0.457 | 3.153 | $\begin{gathered} \text { bO-C(16), bC- } \\ C(48) \end{gathered}$ |
| 52. |  | 538w | 464 | 521 | 3.4745 | 0.21 | 0.4804 | 3.7812 | 465 | 518 | 3.8813 | 0.27 | 0.4831 | 3.7836 | $\begin{gathered} \text { bO-C(20),vC- } \\ \mathrm{O}(16), \omega \mathrm{C}-\mathrm{C}(13) \end{gathered}$ |
| 53. | 529w |  | 414 | 461 | 0.6312 | 1.91 | 0.6299 | 0.2218 | 413 | 459 | 0.5864 | 1.9 | 0.6167 | 6.1226 | $\begin{aligned} & \text { bC-O(32), bC- } \\ & \mathrm{C}(12), \mathrm{vC}-\mathrm{C}(10) \end{aligned}$ |


| 54. |  | 512w | 390 | 455 | 4.4975 | 2.79 | 0.6395 | 7.1072 | 390 | 451 | 3.9782 | 2.39 | 0.6035 | 6.7057 | $\begin{gathered} \operatorname{tRtring}(57), \\ \text { tRasymd(17), } \omega \mathrm{C}- \\ \mathrm{C}(08) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 55. | 454vw |  | 366 | 419 | 3.1834 | 1.48 | 0.2948 | 3.7213 | 365 | 415 | 3.881 | 1.44 | 0.2904 | 3.6948 | $\begin{gathered} \omega \mathrm{C}-\mathrm{C}(22), \mathrm{bO}- \\ \mathrm{H}(16) \end{gathered}$ |
| 56. | 447vw |  | 340 | 303 | 3.1008 | 4.34 | 0.375 | 5.4994 | 340 | 299 | 2.7744 | 4.36 | 0.3778 | 5.5446 | $\begin{gathered} \omega \mathrm{C}-\mathrm{C}(17), \omega \mathrm{O}-\mathrm{C} \\ \text { (27) } \end{gathered}$ |
| 57. | 432s |  | 303 | 279 | 5.2793 | 1.91 | 0.2189 | 4.0336 | 301 | 275 | 5.6777 | 1.61 | 0.2163 | 4.05 | $\begin{gathered} \omega \mathrm{C}=\mathrm{O}(20), \mathrm{vC}- \\ \mathrm{O}(16), \omega \mathrm{C}-\mathrm{C}(13) \end{gathered}$ |
| 58. | 424w |  | 230 | 253 | 2.6287 | 0.03 | 0.175 | 5.5985 | 227 | 249 | 2.4486 | 0.23 | 0.1736 | 5.6761 | $\begin{gathered} \omega \mathrm{C}-\mathrm{C}(24), \mathrm{bC}- \\ \mathrm{O}(10), \omega \mathrm{C}-\mathrm{C}(20) \end{gathered}$ |
| 59. | 418w |  | 210 | 241 | 1.4552 | 1.64 | 0.0945 | 3.6221 | 207 | 237 | 1.6139 | 1.67 | 0.0919 | 3.6246 | $\begin{gathered} \omega \mathrm{C}=\mathrm{O}(18), \mathrm{vC}- \\ \mathrm{O}(16), \omega \mathrm{C}-\mathrm{C}(13) \end{gathered}$ |
| 60. |  | 412w | 182 | 216 | 0.0686 | 0.57 | 0.022 | 1.1212 | 171 | 212 | 0.0496 | 0.97 | 0.0192 | 1.114 | $\begin{gathered} \omega \mathrm{O}-\mathrm{C}(19), v \mathrm{C}- \\ \mathrm{O}(16), \omega \mathrm{C}-\mathrm{C}(13) \end{gathered}$ |
| 61. |  | 294ms | 139 | 137 | 0.4512 | 0.54 | 0.0178 | 1.5609 | 137 | 134 | 0.7568 | 0.58 | 0.0204 | 1.8452 | $\begin{gathered} \omega \mathrm{O}-\mathrm{C}(20), \nu \mathrm{C}- \\ \mathrm{O}(08), \omega \mathrm{C}-\mathrm{C}(13) \end{gathered}$ |
| 62. |  | 230w | 132 | 112 | 0.7432 | 1.04 | 0.0192 | 1.8516 | 129 | 109 | 0.6526 | 1.34 | 0.0157 | 1.5815 | $\begin{gathered} \omega \mathrm{O}-\mathrm{C}(20), v \mathrm{C}- \\ \mathrm{O}(16), \omega \mathrm{C}-\mathrm{C}(13) \end{gathered}$ |
| 63. |  | 142w | 73 | 84 | 0.7816 | 1.07 | 0.0131 | 4.1042 | 73 | 81 | 0.714 | 1.97 | 0.0127 | 3.9753 | tRsymd(51), tRtring(33), bC$\mathrm{C}(10)$ |
| 64. |  | 127s | 57 | 76 | 6.1756 | 1.83 | 0.0091 | 4.6943 | 56 | 72 | 5.985 | 1.87 | 0.0084 | 4.4099 | tRasymd(61), tRsymd(27), vC C(11) |


| 65. | 101vs | 42 | 63 | 1.3241 | 2.16 | 0.0057 | 5.4214 | 41 | 61 | 1.7806 | 2.19 | 0.0066 | 6.6502 | $\tau \mathrm{CH}_{3}(69)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 66. | 75vs | 31 | 57 | 1.9438 | 0.75 | 0.0048 | 7.9395 | 34 | 54 | 1.2109 | 0.89 | 0.0051 | 7.2246 | $\tau \mathrm{CH}_{3}$ (55) |

Abbreviations: vs: very strong; ms: medium strong; s: strong; w: weak; mw: medium weak; vw: very weak; ips: inplane stretching; ops: out-of-plane stretching; ss: symmetric stretching; ipb: in-plane bending; opb:
IR out-of-plane bending; sb: stretching bending; $\tau$ : twisting.

## Methyl Group Vibrations

The 3A2MBA molecule consists of two $\mathrm{CH}_{3}$ groups, one unit in benzene ring and another in the acetoxy group, which lies in the tertiary position of the molecule. For the assignments of $\mathrm{CH}_{3}$ group frequencies, nine fundamentals can be associated to each $\mathrm{CH}_{3}$ group[53] namely $\mathrm{CH}_{3}$ ss-symmetric stretch; $\mathrm{CH}_{3}$ ips-in-plane stretch (i.e. in-plane hydrogen stretching modes); $\mathrm{CH}_{3} \mathrm{ipb}$ - in-plane bending (i.e. in-plane hydrogen deformation modes); $\mathrm{CH}_{3}$ sb-symmetric bending; $\mathrm{CH}_{3} \mathrm{ipr}$-in-plane rocking; $\mathrm{CH}_{3}$ ops-out-of-plane stretching; $\mathrm{CH}_{3}$ opb-out-ofplane bending; $\mathrm{CH}_{3} \mathrm{Opr}$-out-of-plane rocking and $\tau \mathrm{CH}_{3}$-torsion. The $\mathrm{C}-\mathrm{H}$ stretching in $\mathrm{CH}_{3}$ occurs at lower frequencies than those of aromatic ring. Moreover, the asymmetric stretch is usually at higher wavenumber than the symmetric stretch. Methyl group vibrations are generally referred to as electron-donating substituent in the aromatic rings system. Acetoxy group comprises a methyl and an ester group. The hydrogen in the methyl group acts as an electron donor, whereas oxygen present in the ester is an electron acceptor. The anti-symmetric and symmetric stretching vibrations are in the ranges 3050-2950 and 2970-2860 $\mathrm{cm}^{-1}$ [54]. In the present study, two $\mathrm{CH}_{3}$ anti symmetric stretching vibrations are observed at 2999 and $2941 \mathrm{~cm}^{-1}$ in FT-Raman and FT-IR respectively. Two $\mathrm{CH}_{3}$ symmetric stretching modes are observed as weak bands in FT-IR and FT-Raman spectra at 2894 and $2920 \mathrm{~cm}^{-1}$ of title molecule and the corresponding two $\mathrm{CH}_{3} \mathrm{ips}$, in-plane stretching modes are observed in FT-IR spectrum at 2882 and $2847 \mathrm{~cm}^{-1}$. The $\mathrm{CH}_{3}$ asymmetric and symmetric stretching vibrations predicted theoretically by B3LYP/6-311++G (d,p) method show excellent agreement with the experimental values. It shows above $90 \%$ of TED contribution suggesting that it is pure stretching mode. For methyl substituted benzene derivatives the antisymmetric and symmetric deformation vibrations of methyl group normally appear in the regions $1480-1410$ and $1385-1250 \mathrm{~cm}^{-1}$ [54,55] respectively. In the FT-IR and FTRaman spectra, the bands occurred at $1425,1364,1317,1270$ and $1412,1300 \mathrm{~cm}^{-1}$ are due to antisymmetric and symmetric deformations. The computed values are predicted by B3LYP/6-31+G (d,p) and 6-311++G (d,p) method show good agreement with literature as well as experimental observations. The rocking vibrations of $\mathrm{CH}_{3}$ modes usually appear in the region $1070-826 \mathrm{~cm}^{-1}$ [56]. The Strong and medium weak bands observed at 936 and $882 \mathrm{~cm}^{-1}$ in FT-IR Spectrum are assigned to two $\mathrm{CH}_{3} \mathrm{ipr}-$ in-plane rocking vibrations and its counterpart, $\mathrm{CH}_{3}$ opr-out-of-plane rocking vibrations were observed at 871 and $835 \mathrm{~cm}^{-1}$ in FT-Raman and FTIR spectra. The experimental and theoretically wavenumbers are predicted of $\mathrm{CH}_{3}$ rocking vibrations show good correlation with the literature data of the title molecule. The $\mathrm{CH}_{3}$ torsional modes of vibrations were assigned in the lower region of the Raman spectrum and the values are in line with the theoretical values.

## Molecular Orbital Studies

The most widely used theory by chemists is the molecular orbital (MO) theory. It is important that Ionization Potential (IP), Electron Affinity (EA), Electrophilicity Index ( $\omega$ ), Chemical Potential ( $\mu$ ), Electronegativity ( $\chi$ ) and Hardness $(\eta)$ be put into a MO framework. These are readily being done within the limitations of Koopman's theorem the orbital energies of the frontier orbital are given by Figure 4 show an orbital energy diagram for 3A2MBA molecule. Only the HOMO and LUMO orbitals are shown. On the basis of a fully optimized ground state structure, the DFT/B3LYP/6-31+G(d,p) and $6-311++G(d, p)$ calculations predicts one intense electronic transition from the ground to the first excited state and is mainly described by one electron excitation from the highest occupied molecular orbital (HOMO) to the lowest unoccupied molecular orbital (LUMO). We focus on the HOMO and LUMO energies in order to determine interesting molecular/atomic properties and chemical quantities. In simple molecular orbital theory approaches, the HOMO energy (HOMO) is related to the IP by Koopmans's theorem and the LUMO energy (LUMO) has been used to estimate the electron affinity (EA) [57]. If HOMO and LUMO, then the average value of the HOMO and LUMO energies is related to the electronegativity $(\chi)$ defined by Mulliken [58] with $\chi=(\mathrm{IP}+\mathrm{A}) / 2$. In addition, the HOMOLUMO gap is related to the hardness ( $\eta$ ) [59-62], and also Parr et al. [57] defined global electrophilicity as a quantitative intrinsic numerical value and suggested the term electrophilicity index ( $\omega$ ), a new global reactivity descriptor of molecule, as $=\mu^{2} / 2$ where $\mu$ is the chemical potential takes the average value [61] i.e., $\mu=$ (IP $+E A) / 2$. In general, the electrophiles are electron lovers or electron deficient and hence prefer to accept electrons and form bonds with nucleophiles. Thus electrophilicity is useful structural depictor of reactivity and is frequently used in the analysis of the chemical reactivity of molecule [62]. The title molecule with low potential is a good electrophile; an extremely hard molecule has feeble electron acceptability. Consequently, a measure of molecular electrophilicity depends on both the chemical potential and the chemical hardness. Finally, the calculated HOMO-LUMO gaps are also closer to the first electronic excitation energies. The electronegativity and hardness are of course used extensively to make predictions about chemical behavior and these are used to explain aromaticity in organic compounds [62]. A hard molecule has a large HOMO-LUMO gap and a soft molecule has a small HOMO-LUMO gap. In quantum theory, changes in the electron density of a chemical system result from the mixing of suitable excited-state wave function with the ground state wave function. The mixing coefficient is inversely proportional to the excitation energy between the ground and excited state. A small HOMO-LUMO gap automatically means small excitation energies to the manifold of
excited states. Therefore, soft molecule, with a small gap, will have their electron density changed more easily than a hard molecule. In terms of chemical reactivity, we can conclude that soft molecule will be more reactive than hard molecule. The HOMO is located over methyl group atoms and acetoxy group atoms, and the HOMOLUMO transition implies an electron density transfer to the carboxylic group. Moreover, these orbitals significantly overlap in their position for 3A2MBA is shown in Figure 4. The HOMO-LUMO energy gap of 3A2MBA were calculated at the B3LYP/6-31+G(d,p) and $6-311++G(d, p)$ levels, which reveals that the energy gap reflects the chemical activity of the molecule. The LUMO as an electron acceptor represents the ability to obtain an electron and HOMO represents the ability to donate an electron. The energy gap between the HOMO and the LUMO molecular orbitals is a critical parameter in determining molecular electrical transport properties because it is a measure of electron conductivity. The energy gap values of HOMO-LUMO are -4.09151 and -5.23710 eV of 3A2MBA, respectively. Lower value in the HOMO and LUMO energy gap explains the eventual charge transfer interactions taking place within the molecule. The calculated chemical hardness, softness, chemical potential, electrophilicity, ionization potential, electron affinity and electronegativity are presented in Table 3.

## Natural Atomic Orbitals

The occupancies and energies of bonding molecular orbitals of the title molecule is predicted at B3LYP/6$311++G(d, p)$ level of theory and presented in Table 4. The variations in occupancies and energies of 3A2MBA directly give the evidence for the delocalization of charge upon substitution and this leads to the variation of bond lengths of the title molecule as shown in Table 1.


Figure 4: The atomic orbital compositions of the frontier molecular orbital for 3A2MBA along with B3LYP/ 6-311++G (d,p) method basis set
Table 3: Calculated molecular properties of 3A2MBA from orbital energies

| Chemical parameters( in eV) | Values of B3LYP |  |
| :---: | :---: | :---: |
|  | $\mathbf{6 - 3 1 + G ( d , p )}$ | $\mathbf{6 - 3 1 1 + G ( d , p )}$ |
| E $_{\text {Homo }}$ | -9.4195 | -7.31034 |
| ELUMO | -5.32799 | -2.07323 |
| Energy gap | -4.09151 | -5.2371 |
| Ionization potential (I) | 9.4195 | 7.31034 |
| Electron affinity (A) | 5.32799 | 2.07323 |
| Electronegativity(c) | 7.37374 | 4.69179 |
| Global hardness (h) | 2.04575 | 2.61856 |
| Chemical potential $(\mu)$ | -7.37374 | -4.691789 |
| Global Electrophilicity $(\omega)$ | 13.28895 | 4.20334 |
| Softness(S) | 361.94993 | 282.7743 |

Table 4: Comparison of occupancies and energies of bonding molecular orbitals for 3A2MBA using B3LYP/6-311++G (d,p) method

| Atomic orbitals $^{\mathbf{a}}$ | Occupancies(e) | Energies(a.u) |
| :---: | :---: | :---: |
| $\mathrm{BD}(\mathrm{C} 1-\mathrm{C} 2)$ | 1.97273 | -0.7197 |
| $\mathrm{BD}(\mathrm{C} 1-\mathrm{C} 6)$ | 1.97246 | -0.71329 |
| $\mathrm{BD}(\mathrm{C} 1-\mathrm{C} 7)$ | 1.97302 | -0.65816 |
| $\mathrm{BD}(\mathrm{C} 2-\mathrm{C} 3)$ | 1.95601 | -0.71169 |
| $\mathrm{BD}(\mathrm{C} 2-\mathrm{C} 11)$ | 1.97904 | -0.61442 |
| $\mathrm{BD}(\mathrm{C} 3-\mathrm{C} 4)$ | 1.96243 | -0.70936 |
| $\mathrm{BD}(\mathrm{C} 3-\mathrm{O} 15)$ | 1.98627 | -1.36283 |
| $\mathrm{BD}(\mathrm{C} 4-\mathrm{C} 5)$ | 1.98051 | -0.70994 |
| $\mathrm{BD}(\mathrm{C} 4-\mathrm{H} 22)$ | 1.97952 | -0.54334 |
| $\mathrm{BD}(\mathrm{C} 5-\mathrm{C} 6)$ | 1.97749 | -0.70788 |
| $\mathrm{BD}(\mathrm{C} 5-\mathrm{H} 23)$ | 1.98053 | -0.53596 |
| $\mathrm{BD}(\mathrm{C} 6-\mathrm{H} 24)$ | 1.9783 | -0.52809 |
| $\mathrm{BD}(\mathrm{C} 7-\mathrm{O})$ | 1.99618 | -1.06411 |
| $\mathrm{BD}(\mathrm{C} 7-\mathrm{O} 9)$ | 1.99549 | -0.94929 |
| $\mathrm{BD}(\mathrm{O} 9-\mathrm{H} 10)$ | 1.98673 | -0.77435 |


| $\mathrm{BD}(\mathrm{C} 11-\mathrm{H} 12)$ | 1.97635 | -0.51312 |
| :---: | :---: | :---: |
| $\mathrm{BD}(\mathrm{C} 11-\mathrm{H} 13)$ | 1.97711 | -0.51607 |
| $\mathrm{BD}(\mathrm{C} 11-\mathrm{H} 14)$ | 1.98875 | -0.81925 |
| $\mathrm{BD}(\mathrm{O} 15-\mathrm{C} 16)$ | 1.93138 | -1.0999 |
| $\mathrm{BD}(\mathrm{C} 16-\mathrm{O} 17)$ | 1.99622 | -0.67934 |
| $\mathrm{BD}(\mathrm{C} 16-\mathrm{C} 18)$ | 1.9848 | -0.67934 |
| $\mathrm{BD}(\mathrm{C} 18-\mathrm{H} 19)$ | 1.96713 | -0.55508 |
| $\mathrm{BD}(\mathrm{C} 18-\mathrm{H} 20)$ | 1.97168 | -0.55207 |
| $\mathrm{BD}(\mathrm{C} 18-\mathrm{H} 21)$ | 1.98172 | -0.54941 |

## Natural Population Analysis

The natural population analysis [63] performed on the electronic structure of 3A2MBA clearly describes the distribution of electrons in various sub-shells of their atomic orbitals. The accumulation of charges on the individual atom and the accumulation of electrons in the core, valence and Rydberg sub-shells are also presented in Table 5. In the case of 3A2MBA, the most electropositive charge of 0.80034 e is accumulated on C7 and 0.779092 e is accumulated on C16 atom. According to an electrostatic point of view of the molecule, these electropositive atoms have a tendency to accept an electron. Whereas, the most electronegative atoms such as; O15 ( $-0.46472 \mathrm{e})$, O17 ( -0.52325 e ), O8 ( -0.60688 e ) and O9 ( -0.70951 e ) have a tendency to donate an electron. The substituent effect may cause a change in the magnitude of these charges observed in 3A2MBA. In view of title molecule, the electronegative atoms $\mathrm{O} 8, \mathrm{O} 9, \mathrm{O} 15$ and O 17 are of variations in magnitude of natural charges on C 7 and C 16 of electropositive atoms. Further, the natural population analysis showed that 102 electrons in the 3A2MBA molecule are distributed on the sub-shells as follows:
Core: 27.98971 ( $99.9633 \%$ of 28); Valence: 73.67036 ( $99.5545 \%$ of 74); Rydberg: 0.33993 ( $0.3333 \%$ of 102).
Table 5: Accumulation of natural charges, population of electrons in core, valence and Rydberg orbitals of 3A2MBA using
B3LYP/6-311++G (d,p) method

| Atom | Charges (e) | Natural population(e) |  |  | Total (e) |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Core | Valence | Rydberg |  |
| C1 | -0.15182 | 1.99877 | 4.13422 | 0.01884 | 6.15182 |
| C2 | -0.02726 | 1.99896 | 4.00977 | 0.01853 | 6.02726 |
| C3 | 0.28526 | 1.99884 | 3.67513 | 0.04076 | 5.71474 |
| C4 | -0.23409 | 1.99914 | 4.21801 | 0.01695 | 6.23409 |
| C5 | -0.19309 | 1.99921 | 4.17679 | 0.01708 | 6.19309 |
| C6 | -0.18491 | 1.99909 | 4.1679 | 0.01792 | 6.18491 |
| C7 | 0.80034 | 1.99944 | 3.15158 | 0.04864 | 5.19966 |
| O8 | -0.60688 | 1.99976 | 6.59455 | 0.01257 | 8.60688 |
| O9 | -0.70951 | 1.99975 | 6.69587 | 0.01389 | 8.70951 |
| H10 | 0.49756 | 0 | 0.49849 | 0.00396 | 0.50244 |
| C11 | -0.59846 | 1.99927 | 4.58713 | 0.01207 | 6.59846 |
| H12 | 0.23382 | 0 | 0.7648 | 0.00138 | 0.76618 |
| H13 | 0.22003 | 0 | 0.77845 | 0.00152 | 0.77997 |
| H14 | 0.19643 | 0 | 0.80208 | 0.00149 | 0.80357 |
| O15 | -0.46472 | 1.99921 | 6.43312 | 0.03239 | 8.46472 |
| C16 | 0.77909 | 1.99936 | 3.16952 | 0.05204 | 5.22091 |
| O17 | -0.52325 | 1.99975 | 6.51103 | 0.01248 | 8.52325 |
| C18 | -0.68548 | 1.99917 | 4.67778 | 0.00853 | 6.68548 |
| H19 | 0.23819 | 0 | 0.76043 | 0.00138 | 0.76181 |
| H20 | 0.24442 | 0 | 0.75409 | 0.00149 | 0.75558 |
| H21 | 0.23974 | 0 | 0.75918 | 0.00108 | 0.76026 |
| H22 | 0.2102 | 0 | 0.78808 | 0.00173 | 0.7898 |
| H23 | 0.20855 | 0 | 0.79003 | 0.00142 | 0.79145 |
| H24 | 0.22584 | 0 | 0.77236 | 0.0018 | 0.77416 |

## Natural Bond Orbital Analysis

The information about the charge delocalization and the chemical reactivity is being helpful to the drug designers to design new type of drugs. The natural bond orbital (NBO) is an effective tool for the elucidation of residual resonance delocalization effects of a molecule and it also illustrates the deciphering of the molecular wave function in terms of Lewis structures, charge, bond order, bond type, hybridization, resonance, donoracceptor interactions, etc. [64]. Therefore, the NBO analysis of certain pharmaceutical compounds has been performed by various spectroscopists [65-70]. The natural bond orbital (NBO) analysis of 3A2MBA molecule is being performed at DFT/B3LYP/6-311+G level to estimate the delocalization patterns of electron density (ED) from the principle occupied Lewis-type (bond or lone pair) orbitals to unoccupied non-Lewis (anti-bonding or Rydberg) orbitals. Table 6 lists the occupancies and energies of most interacting NBO's along with their percentage of hybrid atomic orbital [71] contribution. The percentage of hybrid atomic orbitals of oxygen lone pair atoms O8, O9, O15 and O17 of 3A2MBA showed that they are partially contributed to both $s$-type and $p$ -
type sub-shells. In contrast, all the anti-bonding orbitals of 3A2MBA are mainly contributed to p-type sub-shell. The interactions result in a loss of occupancy from the localized NBO of the idealized Lewis structure into an empty Non-Lewis orbital. For each donor $(i)$ and acceptor $(j)$, the stabilization energy, $E^{(2)}$ associated with the delocalization $i j$ is estimated as:

$$
\begin{equation*}
\mathrm{E}^{(2)}=\Delta \mathrm{E}_{\mathrm{ij}}=\mathrm{q}_{i} \frac{\mathrm{~F}(\mathrm{i}, \mathrm{j})}{\varepsilon_{\mathrm{j}}-\varepsilon_{\mathrm{i}}} \ldots \tag{2}
\end{equation*}
$$

Where $q_{i}$ is the donor orbital occupancy, $\varepsilon_{i}$ and $\varepsilon_{j}$ are diagonal elements and $F(i, j)$ is the off diagonal NBO Fock matrix element. The NBO analysis provides an efficient method for studying intra-and intermolecular bonding hyperconjugative interactions of $\pi$ electrons and also provides a convenient basis for investigating charge transfer or conjugative interactions in molecular systems [72]. The transferring of energy during the intermolecular interactions between LP (2) O (9) and $\pi^{*}$ (C7-O8) bond is more, which is about $44.79 \mathrm{kj} / \mathrm{mol}$ for 3A2MBA. The maximum energy delocalization takes part in the $\pi^{*} \rightarrow \pi^{*}$ transition. In 3A2MBA, the intra molecular hyper conjugative interaction of the $\sigma(\mathrm{C} 1-\mathrm{C} 2)$ and $\pi(\mathrm{C} 1-\mathrm{C} 2)$ are distributed to $\sigma^{*}(\mathrm{C} 3-\mathrm{O} 15) \rightarrow 4.41 \mathrm{kj} / \mathrm{mol}$ and $\pi^{*}(\mathrm{C} 3-\mathrm{C} 4) \rightarrow 19.15 \mathrm{kj} / \mathrm{mol}, \pi^{*}(\mathrm{C} 5-\mathrm{C} 6) \rightarrow 20.81 \mathrm{kj} / \mathrm{mol}, \pi^{*}(\mathrm{C} 7-\mathrm{O} 8) \rightarrow 20.77 \mathrm{kj} / \mathrm{mol}$. Whereas $\pi^{*}(\mathrm{C} 7-\mathrm{O} 8), \pi^{*}$ (C3-C4) and $\pi^{*}$ (C1-C2), $\pi^{*}$ (C5-C6) bonds are more, which is about 105.33 and $233.77 \mathrm{kj} / \mathrm{mol}$ for 3A2MBA at ground state. The $\mathrm{E}^{(2)}$ values and types of the transition is shown in Table 6.
Table 6: Second order perturbation theory analysis of fock matrix in NBO basis for 3A2MBA using B3LYP/6-311++G(d,p) method

| Donor(i) | Acceptor(j) | E(2) ${ }^{\text {a }}$ | E(j) -E(i) ${ }^{\text {b }}$ | $\mathbf{F}(\mathbf{i}, \mathbf{j})^{\text {c }}$ |
| :---: | :---: | :---: | :---: | :---: |
|  |  | (kj mol ${ }^{-1}$ ) | $\times 10^{3}\left(\mathrm{kj} \mathrm{mol}^{-1}\right)$ | $\times 10^{3}\left(\mathrm{kj} \mathrm{mol}^{-1}\right)$ |
| $\sigma$ (C1-C2) | $\sigma^{*}(\mathrm{C} 3-\mathrm{O} 15)$ | 4.41 | 1.02 | 0.06 |
| $\sigma$ (C3-C4) | $\sigma^{*}$ (C2-C3) | 3.61 | 1.27 | 0.061 |
| $\pi$ (C1-C2) | $\pi^{*}$ (C3-C4) | 19.15 | 0.27 | 0.065 |
|  | $\pi^{*}$ (C5-C6) | 20.81 | 0.28 | 0.07 |
|  | $\pi^{*}$ (C7-08) | 20.77 | 0.26 | 0.067 |
| $\pi$ (C3-C4) | $\pi^{*}$ (C1-C2) | 21.38 | 0.29 | 0.071 |
|  |  | 17.01 | 0.3 | 0.064 |
| $\pi$ (C5-C6) | $\pi^{*}$ (C1-C2) | 18.51 | 0.28 | 0.064 |
|  | $\pi^{*}$ (C3-C4) | 22.71 | 0.27 | 0.07 |
| LP(2) O9 | $\pi^{*}$ (C7-08) | 44.79 | 0.34 | 0.113 |
| LP(2) O15 | $\pi^{*}(\mathrm{C} 16-\mathrm{O} 17)$ | 40.18 | 0.34 | 0.104 |
| LP(2) O17 | $\pi^{*}$ (O15-C16) | 37.13 | 0.59 | 0.133 |
|  | $\pi^{*}$ (C16-C18) | 18.33 | 0.62 | 0.098 |
| $\pi^{*}$ (C3-C4) | $\pi^{*}$ (C5-C6) | 233.77 | 0.01 | 0.081 |
| $\pi^{*}$ (C7-O8) | $\pi^{*}$ (C1-C2) | 105.33 | 0.02 | 0.072 |

${ }^{\mathrm{a}} \mathrm{E}^{(2)}$ means energy of hyper conjugative interactions (stabilization energy); ${ }^{\mathrm{b}}$ Energy difference between donor and acceptor i and j NBO orbitals; ${ }^{\mathrm{c}} \mathrm{F}(\mathrm{i}, \mathrm{j})$ is the Fock matrix element between i and j NBO orbitals

## Thermodynamic Properties

The computation of vibrational frequencies simultaneously leads to prediction of several thermodynamic properties too. The equations utilized for computing thermochemical data in Gaussian programs are derived from statistical thermodynamics with the two key ideas, the Boltzmann distribution and the partition function. Several calculated thermodynamic parameters are presented in prediction in Table 7.

Table 7: The calculated thermodynamic parameters of 3A2MBA

| Thermodynamic properties | B3LYP |  |
| :---: | :---: | :---: |
|  | $\mathbf{6 - 3 1 + G ( d , p )}$ | $\mathbf{6 - 3 1 1 + + G ( d , p )}$ |
| SCF energy | -691.0405882 | -688.0303882 |
| Zero-point vibrational energy (kcal/mol) |  |  |
| 116.85996 |  |  |
| Rotational temperature (K) | 0.06473 | 0.06472 |
| A | 0.02375 | 0.02073 |
| B | 0.01955 | 0.01725 |
| C |  |  |
| Rotational constants (GHz) | 1.34867 | 1.34853 |
| A | 0.4948 | 0.4319 |
| B | 0.40741 | 0.35942 |
| C | 0.0574 | 0.0574 |
| Dipole moment (Debye) | -0.5519 | -0.5519 |
| $\mu_{\mathrm{x}}$ | 1.1273 | 1.1273 |
| $\mu_{\mathrm{y}}$ |  | 1.2565 |
| $\mu_{\mathrm{z}}$ | 100.511 | 117.285 |
| $\mu_{\text {Total }}$ | 41.694 | 41.694 |


| Rotational | 31.448 | 31.707 |
| :---: | :---: | :---: |
| Vibrational | 27.369 | 43.884 |
| Total heat capacity at constant volume ( $\mathrm{cal} / \mathrm{mol} / \mathrm{K}$ ) | 42.776 | 49.119 |
| Translational | 2.981 | 2.981 |
| Rotational | 2.981 | 2.981 |
| Vibrational | 36.814 | 43.157 |
| Total thermal energy (kcal/mol) | 123.502 | 123.782 |
| Translational | 0.889 | 0.889 |
| Rotational | 0.889 | 0.889 |
| Vibrational | 121.724 | 122.005 |
| Thermal corrections (Hartree/Particle) |  |  |
| Zero-point | 0.186228 | 0.183897 |
| Energy | 0.196812 | 0.197259 |
| Enthalpy | 0.197756 | 0.198204 |
| Gibbs free energy | 0.15 | 0.142478 |
| Sum of thermal corrections and energy (Hartree) |  |  |
| Zero-point energies | -687.84416 | -688.021722 |
| Thermal energies | -687.833576 | -688.008359 |
| Thermal enthalpy | -687.832632 | -688.007415 |
| Thermal free energy | -687.880388 | -688.063141 |

Scale factors have been recommended for an accurate determining the zero-point vibration energies (ZPVE), rotational constants, thermal energy and the entropy of the 3A2MBA was directly obtained from the output of Gaussian programs in the optimization and vibration calculations. These attributes includes such as enthalpy $\left(\mathrm{H}_{\mathrm{m}}^{\mathrm{o}}\right)$, entropy $\left(\mathrm{S}_{\mathrm{m}}^{\mathrm{o}}\right)$, heat capacity at constant volume $\left(\mathrm{C}_{\mathrm{v}}\right)$ and pressure $\left(\mathrm{C}_{\mathrm{p}}\right)$ which are contributed by translational, rotational and vibrational partition functions [73]. In addition, SCF energy, zero point vibrational energy, total thermal energy, rotational constants, thermal corrections and dipole moment in all the three dimensions of space were calculated and given in Table 7 for 3A2MBA under B3LYP/6-31+G(d,p) and B3LYP/6-311++G(d,p) model chemistry. All these thermodynamic data provide helpful information on further study of this title molecule, while using as reactant to take part in a new reaction, chemical physics analysis as well as to propose organic reaction mechanisms. These standard thermodynamic functions can be used as reference thermodynamic values to calculate changes which take place in entropy, Gibbs free energy, internal energy, heat capacity and enthalpy of 3A2MBA in its organic chemical transformation [74,75].

## CONCLUSION

A complete vibrational analysis of 3-acetoxy-2-methylbenzoic acid was performed based on calculation at the B3LYP/6-31+G(d,p) and $6-311++G(d, p)$ levels. The theoretical results were compared with the experimental vibrations. Scaling factors results are in good agreement with the experimental values. The computed geometrical parameters are in good agreement with the literature values of this compound. The influences of acetoxy and methyl groups to the vibrational frequencies were also discussed. Natural atomic orbitals, natural population analysis and natural bond orbital analysis are determined. Other electronic structure properties such as chemical hardness, chemical potential, electrophilicity, ionization potential, electron affinity, electronegativity and HOMO-LUMO analysis of the molecule have been calculated. The entire calculated quantum chemical investigations can lay a strong foundation for 3A2MBA to extend its application in organic chemical transformation, material sciences and to model potential bioactive targets.

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